SIMULTANEOUS DETERMINATION OF THE THERMAL CONDUCTIVITY AND DIFFUSIVITY OF METALS USING AN ELECTRONIC HEATER

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A nonstationary method allowing simultaneous determination of the thermal conductivity and diffusivity of metals using an electronic heater is described.

Because of the need to investigate thermophysical properties over a wide temperature range, it is very important to develop nonstationary methods which allow determination in a single experiment of the temperature dependence of one thermophysical property or another. Naturally, methods which afford the maximum information about the properties of the specimen in a single experiment are the most attractive. Another real consideration is the need to devise methods that use a convenient means of heating the specimen and generating a constant heat flux through it, i.e., an electronic heater.

These considerations have led to a new method of simultaneous determination of thermal conductivity and diffusivity of metals, based on the solution of the unsteady heat conduction problem with boundary conditions of the second kind for a semi-infinite rod.

A semi-infinite rod with uniform temperature distribution at time zero is heated by a constant heat flux of value $q_e = \text{const.}$ The temperature varies in one direction. The equation of heat propagation in this case has the form

$$\frac{\partial t(x, \tau)}{\partial \tau} = a \frac{\partial^2(x, \tau)}{\partial x^2} \quad (\tau > 0, \ 0 < x < \infty), \qquad (1)$$

$$t(x, 0) = t_0 = \text{const},$$
 (2)

$$\lambda \frac{\partial t(0, \tau)}{\partial x} + q_e = 0, \qquad (3)$$

$$t(\infty, \tau) = t_0, \qquad (4)$$

$$\frac{\partial t\left(\infty,\,\tau\right)}{\partial x}=0.$$
 (5)

The temperature distribution along the rod has the form [1]

$$t(x, \tau) - t_0 = \frac{2q_e}{\lambda} \sqrt{a\tau} \text{ ierfc } \frac{x}{2\sqrt{a\tau}}$$
 (6)

Solution (6) is valid for the case when there are no energy losses from the side surfaces of the specimen. When there are such losses, we have, instead of (1),

$$\frac{\partial t}{\partial \tau} + a \frac{\partial^2 t}{\partial x^2} + ct = 0, \qquad (7)$$

and boundary condition (3) must be replaced by

$$\lambda \frac{\partial t(0, \tau)}{\partial x} + q_e - ct = 0, \qquad (3a)$$

where the term ct takes account of all possible energy losses in the given experiment.

Equation (7) may be reduced to (1) by the substitution [2]

$$t = \exp\left(-c\,\tau\right)\,t\,(x,\,\tau),\tag{8}$$

and from (7), (8), and (1) we find

$$t(x, \tau) - t_0 = \frac{2q_e \exp(-c\tau)}{\lambda} \sqrt{a\tau} \text{ ierfc } \frac{x}{2\sqrt{a\tau}}.$$
 (9)

Equation (9) may be written for convenience in the form

$$t(x, \tau) - t_{0} = \frac{2q_{e}\exp\left(-c\tau\right)}{\lambda} \sqrt{a\tau} \left\{ \frac{\exp\left[-\left(x/2\sqrt{a\tau}\right)^{2}\right]}{\sqrt{\pi}} - \frac{x}{2\sqrt{a\tau}} \left(1 - \operatorname{erf}\left(\frac{x}{2\sqrt{a\tau}}\right)\right) \right\}.$$
 (10)

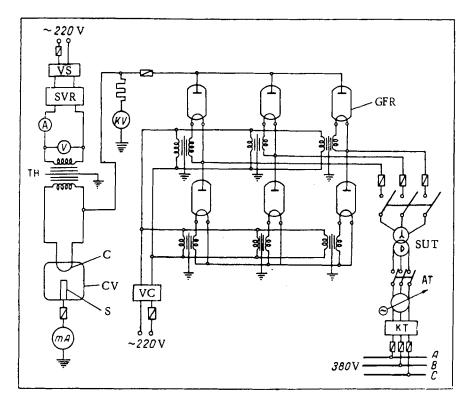
When $\frac{x}{2\sqrt{a\tau}} \le 0.1$, we may neglect the term $\frac{x}{2\sqrt{a\tau}} \times \left(\text{erf } \frac{x}{2\sqrt{a\tau}} \right)$ with an accuracy of 1%, and we may consider the term $\frac{1}{\sqrt{\pi}} \exp \left[-\left(\frac{x}{2\sqrt{a\tau}}\right)^2 \right]$ equal to $1/\sqrt{\pi}$, since the exponent is close to zero. Since the thermal diffusivities of metals are large, it is easy to choose x such that the term $x/2\sqrt{a\tau}$ will always be less than 0.1. Instead of (10) therefore, we may use the simpler expression

$$t(x, \tau) - t_0 = \frac{q_e \exp\left(-c\tau\right)}{\lambda} \left(\frac{2\sqrt{a\tau}}{\sqrt{\pi}} - x\right). \quad (11)$$

By measuring the dependence of temperature on time at three points on the specimen, we may determine from (11) values of diffusivity, conductivity, and the correction coefficient c at the same temperature:

$$\sqrt{a} = \frac{\sqrt{\pi}}{2} \frac{x_2 \exp(-c\tau_2) - x_1 \exp(-c\tau_1)}{\sqrt{\tau_2} \exp(-c\tau_2) - \sqrt{\tau_1} \exp(-c\tau_1)}, \quad (12)$$

$$\lambda = \frac{q_{v} \exp\left(-c\tau_{1}\right)\left(1 - \tau_{2}\right)\left(\tau_{1} - \tau_{2}\right)}{t\left(1 - \exp\left[-c(\tau_{1} - \tau_{2})\right]\sqrt{\tau_{1}/\tau_{2}}\right)},$$
 (13)



Main features of the equipment for simultaneous measurement of thermal conductivity and diffusivity of metals using an electronic heater: VS-voltage stabilizer; SVR-single-phase voltage regulator; TH-heater transformer; C-cathode; VC-vacuum chamber; S-specimen; KV-kilovoltmeter; GFR-gas-filled rectifier; SUT-step-up transformer; ATautotransformer.

$$\exp(c\tau_3)\left(\sqrt{\tau_1}x_2 - \sqrt{\tau_2}x_1\right) =$$

$$= \exp(c\tau_1)\left(\sqrt{\tau_3}x_2 - \sqrt{\tau_2}x_3\right) - \exp(-c\tau_2) \times$$

$$\times (\sqrt{\tau_3}x_1 - \sqrt{\tau_1}x_3). \tag{14}$$

The method outlined allows the thermal conductivity and diffusivity to be measured using various means of creating a constant heat flux along the specimen (lasers, laboratory projector lamps, electronic heaters, etc.). The figure shows one possible variant using an electronic heater. The equipmentworks on the principle of a tube diode, in which the anode (the specimen) is heated by the flux of electrons emitted from the heated cathode. The three-phase high-voltage rectifier is hooked up to the six gasfilled rectifiers in a Larionov bridge circuit. The heater power is measured with a kilovoltmeter and milliameter. Temperature is measured by recording it at three points, with a potentiometric recorder (at low temperatures) or with a pyrometer or photographic method (at high temperatures). The specimen properties are determined by analyzing the $t = f(\tau)$ curves obtained from (14), (12), and (13).

NOTATION

 q_e -heat flux density; t-temperature; x-distance to thermocouple; τ -heating time; t_0 -temperature at time zero; λ -thermal conductivity of material examined; *a*-thermal diffusivity of material; c-coefficient describing the energy losses; x_1 , x_2 , x_3 distance to first, second, and third thermocouples, respectively; τ_1 , τ_2 , τ_3 -times during which specimen is heated to temperature t at points x_1 , x_2 , x_3 , respectively.

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